Empirical Bond-Strength-Bond-Length Curves for Oxides

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Bond-strength-bond-length relationships for bonds between oxygen and H⁺, Li⁺, Be²⁺, B³⁺, Na⁺, Mg²⁺, Al³⁺, Si⁴⁺, P⁵⁺, S⁶⁺, K⁺, Ca²⁺, Sc³⁺, Ti⁴⁺, V⁵⁺, Cr⁶⁺, Mn²⁺, Fe³⁺, Fe²⁺, Co²⁺, Cu²⁺, Zn²⁺, Ga³⁺, Ge⁴⁺ and As⁵⁺ have been derived by requiring that the sums of the bond strengths around the cations be equal to their valence in 417 crystals whose structures have been accurately determined. The relationship is of the form $s = (R/R_0)^{-N}$ where s = bond strength, R = bond length and R_0 and N are fitted constants. It is further shown that all ions with an isoelectronic core can be fitted by a single pair of parameters, R_0 and N, that are independent of the ionic character of the bond and the coordination number of the cation. The resulting bond strength have the property that they are directly related to the covalent character of the bond and that their sum around each atom is, on average, within about 5% of its valence. The bond-strength-bond-length curves are particularly useful in accounting for bonding in cases where the coordination is very distorted (*e.g.* Na⁺, Cu²⁺ and V⁵⁺). They can also be used to predict the positions of hydrogen atoms, to analyze for different oxidation states and site occupancies, to calculate ionic radii and to provide an indication of the correctness of crystal structure determinations.

Introduction

The concept of mean bond strength (\bar{s}) was defined by Pauling (1929) as the valence (z) of a cation divided by its coordination (ν). He enunciated the electrostatic valence principle by which the sums ($p = \sum \bar{s}$) of mean bond strengths around the cations and anions are approximately equal to their valence. This principle has been tested frequently by mineralogists, particularly in silicate structures, and is in general only approximately valid with the sum of the bond strengths around the anions deviating by as much as 40% from the valence (Baur, 1970).

In the original formulation of the principle, the variation of bond strength with bond length was determined by the factor $1/\nu$. It is possible, however, to relate bond strength (s) to bond length (R) using analytical expressions (R-s curves). Byström & Wilhelmi (1951) used Pauling's (1947) logarithmic relationship for covalent bonds in the form:

$$R - R_1 = -2k \log n \tag{1}$$

between bond length, R, and bond number n,† to evaluate bond strengths in $(NH_4)_2Cr_2O_7$ and V_2O_5 (R_1 = length of a single bond and k is an arbitrary constant). Evans (1960) revised the value of R_1 for V⁵⁺ to analyze valence sums in 10 vanadates containing V⁴⁺ and V⁵⁺. The agreement of calculated and observed valences was for the most part excellent although the bond-strength sum, p_V , around vanadium in V₂O₅ was 5.47 and thus deviated considerably from the ideal value of 5.0. Zachariasen (1954, 1963) and Zachariasen & Plettinger (1959) used the distances in a number of borates and uranates to prepare empirical tables of H⁺-O, B³⁺-O, and U⁶⁺-O ionic bond strengths vs. bond lengths and these resulted in bond-strength sums which did not differ by more than 0·1 valence unit (v.u.) from the ideal values. Clark, Appleman & Papike (1969) used an equation of the form $R = a + bs + cs^2$ where R = bond distance, s = bond strength, and a, b, c are variable parameters for each ion pair to evaluate the strengths of the bonds Na⁺-O, Mg²⁺-O, Ca²⁺-O, Al³⁺-O, Fe³⁺-O and Si⁴⁺-O in pyroxenes. Perloff (1970) assumed a linear relationship between bond length and bond strength for Mo⁶⁺-O distances in

 $Na_3(CrMoO_6O_{24}H_6)$. $8H_2O$.

Donnay (1969) and Donnay & Allman (1970) have devised a scheme for constructing R-s curves using an equation of the form

$$s = s_0 \left(\frac{\bar{R}}{R}\right)^N \tag{2}$$

where $s = \text{strength}^*$ of a bond of length R; $s_0 = \text{ideal strength}$ of the bond of length \bar{R} , the mean value of bond length in an individual polyhedron, and N = a constant which is different for each cation-anion pair and in some cases for different cation sites. This expression is used for values of $R < \bar{R}$. When $R > \bar{R}$ they assume a linear relation between s and R,

$$s = s_0 \frac{R_{\max} - R}{R_{\max} - \bar{R}} .$$
⁽³⁾

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 $[\]dagger$ Bond number is defined as the number of shared electron pairs per bond.

^{*} Although Donnay & Allman (1970) chose to use the term bond valence, we prefer to continue Pauling's original terminology. We have also changed the symbols in equation (2) to be consistent with those of Zachariasen (1931).

Values of R_{max} were determined for most metal-(O, OH) bonds in the periodic table by extrapolating an empirical R-s curve derived from the ionic radii of Shannon & Prewitt (1969) to s=0. They achieve excellent agreement between the valence and the sums of the bond strengths but they are forced to construct unique curves for each structure to be analyzed and in certain cases for even the same cation in different sites, *e. g.* Ca(1) and Ca(2) in pumpellyite (Allmann & Donnay, 1971).

Baur (1970) uses the deviations from Pauling's electrostatic valence principle to predict which bonds will be longer or shorter than the average in a particular crystal. He relates the bond length (R_{ij}) to the sum (p_j) of the mean bond strengths around the anion (j) using a linear expression of the form

$$\mathbf{R}_{ij} = a_i + b_i p_j$$
.

He has determined the values of the empirical constants a_i and b_i for fifteen M_i -O atom pairs using the bond lengths observed in some 130 structures but he requires different parameters for the same cation in a different coordination. His approach works well provided the range of R_{ij} values is not large, but for very distorted environments such as are sometimes found around V⁵⁺, his method does not work satisfactorily (Gopal, 1972).

We have derived a set of empirical bond-strengthbond-length* curves based on the equation

$$s = s_0 \left(\frac{R}{R_0}\right)^{-N} \tag{4}$$

by fitting the constants, s_0 , R_0 and N, so that the sums of the bond strengths around the cations are set as nearly as possible equal to the valence in a large variety of oxide structures.

While the equation we use [equation (4)] is the same as that [equation (2)] used by Donnay & Allman (1970), our approach differs from theirs in a number of significant ways. First, we use equation (4) over the whole range of bond strengths, thereby avoiding the assignment of R_{max} (the maximum bonding distance) and the coordination number which they need in order to calculate \bar{R} . Secondly we use the same curve for all bonds between two atomic species wherever they occur. Thirdly, we obtain the values of s_0 , R_0 and N in a different manner, using the information available in a large number of crystal structures rather than tailoring the values to the specific structures under consideration.

Procedure

The bond strength (s) can be calculated from the bond length (R) using the computer program BOST which contains a variety of functions of the form s = f(R), $[e.g. s = A((R-B)/c)^{-N}$ or $s = 10(A-R)/R_0 + c]$. We have chosen to use the following expression

$$s = s_0 \left(\frac{R}{R_0}\right)^{-N} \tag{5}$$

because it gives a good fit and has the advantage of simplicity. The parameter s_0 (the bond strength associated with a bond of length R_0) is assigned arbitrarily by the user and R_0 and N are constants found by least-square fitting as described below.

An initial set of parameters is chosen and used to calculate the bond strength s_{ij} for the bond between atom *i* and atom *j* using the observed bond length R_{ij} . The sum (p_i) of these bond strengths around the atom *i* is given by

$$p_i = \sum_j s_{ij} \tag{6}$$

and in the ideal case this should be equal to the valence z_i of atom *i*.

The initial set of parameters are then refined by least squares to minimize the function

$$Q = \sum_{i=1}^{m} w_i (z_i - p_i)^2$$
(7)

where *m* is the number of atoms of type *i* used in the calculation and w_i is a weight set equal to $1/\sigma^2(p_i)$, $\sigma(p_i)$ being the standard error in p_i as calculated from $\sigma(R_{ij})$, the standard error in the experimental bond length. Refinement works best if the correlation between the parameters R_0 and N is kept small. This can be done by choosing s_0 so that $R_{ij}/R_0 \sim 1$; R_0 then represents a typical bond length and s_0 the corresponding typical bond strength.

The program calculates a number of measures of the agreement between the valence and the bond strength. In addition to the difference, $(z_i - p_i)$, it calculates the relative difference $(z_i - p_i)/z_i$ and the difference normalized to the standard error $(z_i - p_i)/\sigma(p_i)$, the latter being the quantity which enters into the calculation of Q. The overall agreement for the *m* atoms around which sums have been made is conveniently measured by the r.m.s. relative deviation $D_i = [(\sum (z_i - p_i)^2/z_i^2)/m]^{1/2}$.

Three different procedures were followed in determining the parameters R_0 and N. For cations which exist with several different coordination numbers bondstrength sums around all cations for which reasonably accurate bond lengths were available were calculated. The parameters were then refined by fitting the bond strength sums to the valence of the cation. In cases where the cation normally has only one coordination number this procedure is unsatisfactory since almost any value of N will give an equally good fit. In these cases it was necessary to give much greater weight to

^{*} The concept of bond strength as defined by Pauling (1929) is one derived from an ionic model of chemical bonding. In this paper we show that the concept can be used just as well in situations where the bonding is primarily covalent (see below). The fact that the theory works in a given situation cannot be taken as evidence that the bonding is ionic even though for convenience we have continued to use terms such as 'cation' and 'anion'. The term 'electrostatic bond strength' is used in this paper in a formal sense only and it does not imply an ionic model.

structures with atypical coordination numbers. For example the slopes of the R-s curves for Si⁴⁺-O and Mg²⁺-O were constrained to pass as close as possible to the stishovite and SiP₂O₇ points and to the MgAl₂O₄ and MgV₂O₄ points respectively. Furthermore, since no examples of either tetrahedral Ti⁴⁺O₄ or octahedral P⁵⁺O₆ are known, we were forced to use data from TiCl₄, TiBr₄, and PCl₅. In these cases Pauling radii for Cl⁻ and Br⁻ corrected for CN in a manner similar to that for O⁻² (Shannon & Prewitt, 1969) were subtracted from the average M-X distances. The resulting radii for ¹¹O²⁻ to obtain approximate ¹¹Ti⁴⁺-O and Mg²⁺-O, the R-s curves were constrained to pass through the ¹¹Ti⁴⁺-O and V¹P⁵⁺-O points.

The curves for H^+ , S^{6+} and Cr^{6+} were refined in a different way which makes use of the fact that although these atoms normally occur only with one coordination number the bond lengths can vary considerably. In this case the bond-strength sums around the oxygen atoms as well as around the cations were used. This unfortunately restricted the number of structures that could be used to those for which the bond strengths from the oxygen atoms to other cations in the structure could be calculated but on the other hand it gives a much stronger leverage on the value of *N*. The structures used in the refinements are listed in the Appendix.

Although the interatomic distances published in recent crystal-structure determinations are frequently quoted with standard errors of 0.005 Å or less, it is well known that they are subject to systematic errors larger than this. Principal among these is the error arising from thermal motion which is well understood but difficult to apply because of the need to know how the atomic motions are correlated with each other. In cases of molecular fragments which can be treated as rigid (e. g. SO_4^{2-}) it is possible and customary to make these corrections, but for more weakly bonded groups the correction becomes very uncertain. We have therefore only used distances uncorrected for thermal motion, conscious of the fact that these distances are systematically too low (~ 0.01 to 0.02 Å) and will have a true standard error much larger than the values quoted. A more critical evaluation of the effects of thermal motion on the bond distances would, no doubt, improve the agreement between the bond-strength sums and the valence, but such an evaluation is not possible at present.

The R-s curves are particularly useful in accounting for the variation of bond lengths with cation coordination number, a large effect that has long been recognized (Goldschmidt, Barth, Lunde & Zachariasen, 1926; Pauling, 1927; Zachariasen, 1931). Recently, however, the average distances have been found to depend in some measure on the anion coordination number (Jeffreys & Slaughter, 1963; Slaughter, 1966; Shannon & Prewitt, 1969; Brown & Gibbs, 1969) and the nature of the cations surrounding the oxygen ions (Noll, 1963; Lazarev, 1964; Pant & Cruickshank, 1967; Brown & Gibbs, 1970; Shannon, 1971).

Because this latter dependence is generally not yet well defined, this effect has been ignored. The former dependence can be included by making an *a priori* correction for the anion coordination number using the variation of oxygen radius observed by Shannon & Prewitt (1969). Thus all distances can be 'corrected' to those for two coordinate oxygen by subtraction of the following amounts: 0.01Å (3 coordination), 0.03 Å (4 coordination) and 0.05 Å (6 coordination), *etc.* For comparison the bond strength parameters have been refined using both the 'corrected' and the 'uncorrected' bond distances.

Kálmán (1971) has shown that the average tetrahedral M-O distances (\bar{R}) for atoms in groups IV to VII of the periodic table can be expressed in terms of the valence (equivalent to bond strength) by the simple relation

$$\bar{R} = 0.0113 \ n^2 + R'$$
 (8)

where n=8-z and R' is constant for a given row of the periodic table. This suggests that a similar relation might also hold for the standard distances, R_0 . Rather than use equation (8), we have found that we get a good fit with the function

$$s_0 = \left(\frac{R_0}{R_1}\right)^{-N_1} \tag{9}$$

where R_1 and N_1 are constants for a given row. In most cases the value of N_1 obtained does not differ by more than 1 from the individual N's of the atoms in the row. This allows a single pair of parameters to be used for all the elements in the same row that have isoelectronic ionic cores. Thus one curve of the form

$$s_0 = \left(\frac{R}{R_1}\right)^{-N_1} \tag{10}$$

can be used to calculate bond strengths around Li⁺, Be²⁺, and B³⁺ and another to calculate bond strengths around Na⁺, Mg²⁺, Al³⁺, Si⁴⁺, P⁵⁺, and S⁶⁺ etc. The fact that such a fit can be made implies that the length of the M-O bond of unit strength (R_1) is the same (or nearly the same) for all ions with the same electron core. To the extent to which this is true one can calculate bond strengths for bonds not included in the table (e.g. C-O, N-O, Cl-O; see, for example, Table 15).

Results

Tables 1, 2 and 3 list respectively the refined parameters obtained (1) from 'uncorrected' distances for each individual M–O pair (2) from 'uncorrected' distances for ions having isoelectronic cores and (3) from distances corrected for oxygen coordination for each individual M–O pair. Fig. 1 shows typical R-s curves for Si⁴⁺, Ge⁴⁺ and Ti⁴⁺. The points incidate the values of R that were used in the refinement and are, for convenience, plotted along their respective curve. Fig. 2 shows the curves derived from the universal R-s parameters and gives a good indication of the way in which bond strengths fall off with increasing distance.

Table 1. Individual bond-strength-bond-length parameters for M–O bonds in the expression $s = s_0 (R/R_0)^{-N}$

Cation	<i>s</i> ₀ (v.u.)	R_0 (Å)	Ν
H+	0.5	1.184	2.2
Li+	0.25	1.954	3.9
Be ² +	0.5	1.639	4.3
B^{3+}	1.0	1.375	3.9
Na ⁺	0.166	2.449	5.6
Mg ²⁺	0.333	2.098	5.0
Al ³⁺	0.5	1.909	5.0
Si ⁴ +	1.0	1.625	4.5
P ⁵⁺	1.25	1.534	3.2
S ⁶⁺	1.5	1.466	4·0
K+	0.125	2.833	5.0*
Ca ² +	0.25	2.468	6.0
Sc^{3+}	0.5	2.121	6.0
JC Ti⁴ +	0.666	1.952	4·0
V ⁵⁺	1.25	1.714	5.1
Cr ⁶⁺	1.5	1.648	4·9†
Mn^{2+}	0.333	2·186	5.5
Fe ³⁺	0.5	2·012	5.3
Fe ²⁺	0.333	2·155	5.5
Co^{2+}	0.333	2.133 2.118	5.0
Cu ² +	0.333	2.084	5.3
Zn^{2+}	0.5	1.947	5.0
Ga ³⁺	0.75	1.837	4.8
Ge ⁴⁺	1.0	1.750	5.4
As ⁵⁺		1.681	3·4 4·1
AS '	1.25	1.091	4'1

* Refines to N=4.3 but N=5.0 gives as good agreement. † Refines to N=6.1 but N=4.9 gives better overall agreement.

Table 2. Universal bond-strength-bond-length
parameters for M–O bonds in the expression
$s = (R/R_1)^{-N_1}$

Cations	Number of electrons in cation core	R_1 (Å)	N_1
H+	0	0.86	2.17
$Li^{+}Be^{2}B^{3}$	2	1.378	4.065
Na + Mg ² + Al ³ + Si ⁴ + P ⁵ + S ⁶ +	10	1.622	4.290
$K^{+}Ca^{2+}Sc^{3+}Ti^{4+}V^{5+}Cr^{6+}$	18	1.799	4.483
$Mn^{2+}Fe^{3+}$	23	1.760	5.117
$Zn^{2+}Ga^{3+}Ge^{4+}As^{5+}$	28	1.746	6.020

The agreement between the valence and the bondstrength sums is quite sensitive to the value of $N (\pm 0.2)$ for small ions whose coordination numbers are well defined *e. g.* Be²⁺, B³⁺, Si⁴⁺, Zn²⁺, Ga³⁺ and Ge⁴⁺. In other cases N may be varied within a range of ± 2 without making an appreciable difference to the agreement.* On the other hand the values of R_0 are usually quite well determined, variations of 0.003 Å producing a noticeable change in the observed agreement. Table 3. Bond-strength-bond-length parameters for M–O bonds (corrected for oxygen coordination) used in the expression $s = s_0 (R/R_0)^{-N}$

Cation Li ⁺ Be ²⁺ B ³⁺	s_0 (v.u.) 0·25 0·5 1·0	<i>R</i> ₀ (Å) 1·925 1·611 1·343	N 3·4 3·8 3·7
Na ⁺	0.166	2.421	5.7
Mg ²⁺	0.333	2.076	5.2
Al ³⁺	0.2	1.888	4.6
Si ⁴ +	1.0	1.605	4·0
P ⁵⁺	1.25	1.525	3.2
K+	0.125	2.823	6∙0*
Ca ²⁺	0.25	2.437	5.5
Sc ³⁺	0.2	2.090	6.0
Ti⁴ +	0.666	1.947	4.6
V ⁵⁺	1.25	1.700	4.8
Mn ²⁺	0.333	2.165	5.1
Fe ³⁺	0.2	1.981	5.2
Fe ²⁺	0.333	2.128	4.9
Co ² +	0.333	2.087	5.0
Cu ²⁺	0.333	2.068	5.4
Zn ²⁺	0.2	1.946	5.7
Ga ³⁺	0.75	1.822	5.2
Ge ⁴ +	1.0	1.735	5.3
As ⁵⁺	1.25	1.671	4.9

* Refines to N=4.7 but N=6.0 gives as good agreement.

The relative merits of the universal and individual curves and the individual curves corrected for oxygen coordination can be judged from the agreement indices (r.m.s. relative deviation) summarized in Table 4. Comparison of the individual vs. the universal curves shows better agreement for the individual curves in 15 instances, and worse agreement in only 5 cases.[†] The mean r.m.s. deviation of 4.0% for the individual curves can be compared with 5.4% for the universal curves. Comparison of individual 'corrected' and 'uncorrected' curves shows better agreement for the 'uncorrected' distances in 10 cases and worse agreement in another 12 cases; the mean r.m.s. deviations being 4.6 and 4.4%respectively. Since calculations involving uncorrected distances are simpler and since the individual curves give a better fit than the universal ones, the 'uncorrected' individual curves with parameters listed in Table 1 are to be preferred for most applications.

Discussion

1. *Electrostatic-valence principle*

We have tried to find a functional relationship which uniquely relates bond strength to bond length, regardless of structure type, with the constraint that the sum of the bond strengths equals the valence as proposed

^{*} In certain cases when the refined value of N appears to be unreasonably large or small we have chosen a value different from that obtained by the refinement procedure outlined above. In all these cases, which are explicitly noted in Tables 1 and 3, the agreement is equally good with both values.

[†] The individual curves can always be made to give at least as good a fit as the universal ones since the individual parameters can always be chosen to be the same as the universal parameters. However, the method of fitting the parameters is not one that necessarily minimizes the r.m.s. relative deviation. One would expect the individual parameters to show systematic trends across the periodic table and the universal curves provide a reasonably successful attempt to describe these.

by Pauling (1929). This attempt to evaluate the electrostatic-valence principle follows the general lines adopted by most previous workers but differs from them in that our data, evaluated by least-squares techniques, are more extensive. We have analyzed environments of 884 cations in 417 different structures and have derived R-s curves for M-O bonds for the majority of atoms in the first half of the periodic table in a simple twoparameter form. Table 4 shows that the mean r.m.s. deviation of the valence from the calculated bondstrength sums around cations for individual curves is 4.3%; the agreement is worse for the ions Li⁺, Na⁺, K^+ , Ca^{2+} , Sc^{3+} , and Mn^{2+} . Fig. 3, in which bondstrength sums around most of the cations evaluated in this study are plotted vs. their average bond length. shows that although the mean deviation is about 5%. individual variations of 10-15% are still observed.

A detailed examination of the results for V^{5+} shows that while our model accounts for most of the gross features of the structures significant deviations still remain. Table 5 lists the bond sums around the vanadium atoms used in our calculations. For comparison the sums are given using both the individually fitted curve and the curve for 'cations' with the eighteen-electron core. Although the average deviation from 5.00 is only 4.2%, it is apparent that there are individual deviations of 10–20% in some structures. Some of these deviations can be attributed to uncertainties in our knowledge of the bond lengths (indicated by the errors given for the bond-strength sums) but some have to be attributed to real effects which we cannot as yet explain. That these effects are real is attested, for example, by the bond-

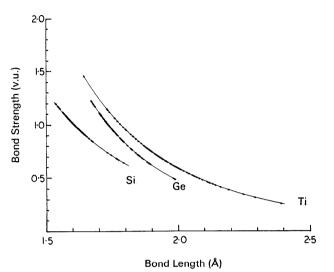


Fig. 1. Individual M–O bond-length-bond-strength curves for M=Si, Ge and Ti. The circles indicate the values of the individual bond lengths used in the calculations. They do not represent the quality of fit between experiment and theory.

Table 4. Statistics concerning the derivation of bond-strength-bond-length parameters

Method 1. Least squares fit R_0 and N using cation sums only.

2. N chosen to give good fit for alypical coordination numbers. R_0 refined by least squares.

3. Least-squares fit of R_0 and N using cation and anion sums.

		Number		R.m.s.	relative deviation	on (%)
	Number of	of bond		Individual	Universal	Individual
	structures	strength		uncorrected	(Table 2)	corrected
Cation	used	sums used	Method	(Table 1)		(Table 3)
H+	9 (3)*	38 (24)*	3	4.6	-	-
Li+	22	29	1	6.4	10.5	5.3
Be ² +	11	13	1	3.4	4.0	3.7
B ^{3 +}	24	44	1	2.6	2.9	4∙0
Na+	31	48	1	7.6	7.5	6.4
Mg ² + Al ³ +	27	39	2	5-2	4.9	5.1
A1 ³ +	25	39	1	4.4	4.8	4.0
Si ^{4 +}	38	50	2	3.1	3.2	2.2
P ^{5 +}	29	45	2	2.1	3.4	2.7
S ⁶⁺	33 (3)*	41 (21)*	3	2.4	. 2.5	_
K+	29	34	1	8.5	10.9	9.0
Ca ²⁺	30	40	1	6.3	8∙4	6.6
Sc ³⁺	8	11	2	7.2	6.1	4.8
Ti4 +	21	41	2	3.3	5-1	4.7
V ⁵ +	31	46	1	4.2	4.1	3.9
Cr ⁶⁺	10 (5)*	22 (50)*	3	3.5	3.5	-
Mn ²⁺	20	33	2	6.9	6.5	6.6
Fe ³⁺	36	50	1	27	2.9	3.2
Fe ²⁺	15	19	1	5.4		4.8
Co ²⁺	18	29	2	6.6	-	3.7
Cu ²⁺	25	40	1	5.8	-	6.2
Zn ²⁺	29	48	1	4.8	4.9	3.9
Ga ³⁺	10	14	1	3.2	6.6	2.8
Ge ⁴⁺	17	28	1	3.7	5.4	4.4
As ⁵⁺	24	34	1	3.1	5.4	3.4

* Figures in parentheses indicate the number used in the refinement [marked with * in the Appendix]. The other figures indicate the number of X-O structures and bonds used in calculating the indices quoted in the last 3 columns.

strength sums consistently found in a number of pyroxene structures where $p_{O(2)}$ is typically 1.75 and $p_{O(3)}$ lies between 2.15 and 2.30. Martin & Donnay (1972) attribute the low value of $p_{O(2)}$ to the presence of OH⁻ ions. However, in the colourless pyroxenes LiVO₃ and NaVO₃, small amounts of hydroxyl ion would be accompanied by reduction of V⁵⁺ to V⁴⁺ which would probably result in darkly coloured compounds. It is likely that the second-nearest neighbours are responsible for some of these deviations but a detailed analysis of their role must await further study.

2. Bond-strength-bond-length curves and implications for chemical bonding

Bond-strength sums calculated using the 10 parameters of the five universal bond-strength curves give al-

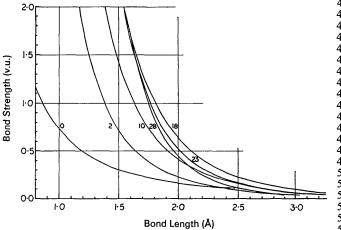


Fig. 2. Universal M–O bond-length–bond-strength curves for isoelectronic series. The numbers associated with each curve indicate the number of electrons on the 'cation' (core electrons).

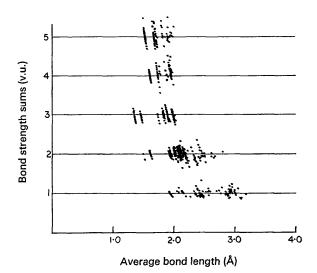


Fig. 3. Experimental bond-length sums vs. average M-O bond length for most of the atoms used in this study.

Table 5. Bond strength sums around V⁵⁺

- v = number of bonds used in calculations given in subsequent columns.
- \bar{R} = average bond length (Å) around V⁵⁺.
- $P_1 =$ bond strength sums calculated with individual R-s curve. $P_1 = \sum_{i} 1.25(R_i/1.714)^{-5\cdot1}$ v.u. (Table 1)

$$P_2$$
 = bond-strength sums calculated with universal $R-s$ curve.
 $P_2 = \sum (R_t/1.799)^{-4.483}$ v.u. (Table 2)

For references see the Appendix.

Figures in parentheses are standard deviations in the last figures quoted.

v	Ŕ	P_1	P_2	Crystal
4	1.715	5.04 (7)	5.01 (6)	$Zn_2V_2O_7$
4	1.706	5.12 (6)	5.08 (5)	YVO4
4	1.721	4.90 (4)	4.89(3)	NdVO₄
4	1.729	4.84 (4)	4.83 (3)	$Mg_3V_2O_8$
4	1.722	4.91 (19)	4.90 (16)	$Co_3V_2O_8$
4	1.740	4.68 (18)	4.69 (16)	$Ni_3V_2O_8$
4	1.719	4.98 (7)	4.96 (6)	$Zn_3V_2O_8$
4	1.718	4.98 (16)	4.95 (14)	$Cd_2V_2O_7$
4	1.720	5.02(5)	4.99 (5)	$Cu_2 v_2 O_7$
4	1.727		4.87 (4)	FeVO₄
4	1.720	4.89 (5)	4.96 (4)	10004
4		4.98(5)		LEVO
4	1.717	4.96(1)	4.94 (1)	Li ₃ VO ₄
4	1.727	4.99 (2)	4.95 (2)	LiVO ₃
4	1·730 1·731	4·92 (7)	4·89 (6) 4·83 (6)	$Co_2V_2O_7$
4	1.707	4·84 (7) 5·11 (4)	5.07 (3)	Ca₂VO₄Cl
4	1.691	5.36(6)	$5 \cdot 28 (5)$	$Na_3VO_4.12H_2O$
4	1.702		$5 \cdot 14(7)$	14a3 V 04.121120
4	1.694	5.20(8)		$Ca_3V_2O_8$
4	1.691	5·31 (10) 5·36 (9)	5·24 (9) 5·29 (8)	$Ca_3 v_2 O_8$
4			4·95 (20)	NH4VO3
4	1·725 1·731	4·99 (23) 4·89 (7)	4.87 (6)	KVO ₃
5	1.819	4·89 (7) 5·09 (7)	5·15 (6)	CaV_2O_6
5	1.848			CaV_2O_6
5	1.808	4.84(22)	4·92 (19) 5·45 (21) [$CaV_2O_6.2H_2O$
5	1.952	5·44 (24) 5·12 (2)	5.12(1)	
5 5 5 5 5	1.883	4.81(1)	4.89(1)	$Mg_2V_2O_7$
5	1.830	4.99 (15)	5.06 (13)	KVO3.H2O
6	2.009	5·29 (11)	5.36 (9)	CsV_3O_8
6	1.906	5.03(2)	5.21(1)	$C_{3}v_{3}O_{8}$
6	1.900	5.03(2) 5.08(3)	$5\cdot21(1)$ $5\cdot22(3)$	
6	1.924		5.17(2)	Ca ₃ V ₁₀ O ₂₈ .17H ₂ O
6	1.910	4·99 (2) 5·05 (2)	$5 \cdot 22 (2)$	
6		1.1	5.22(2) 5.13(4)	
6	1·917 1·913	4.96 (5)		
		5.04(5)	5.20(5)	
6	1.931	4.94 (6)	5.10(5)	$K_2 Z n_2 V_{10} O_{28}$. $16 H_2 O_{28}$
6	1.022	4.85(4)	5.05 (4)	
6 6	1.932	4.94 (5)	5.10(4) J	VPO ₅
	1.947	5.25(5)	5.35 (4)	
6	1.989	5.09(7)	5·20 (6)	V_2O_5
6	1.008	5.27 (36)	5·38 (31) 5·67 (38)	VO(OCH ₃) ₃
6	1.908	5.64(46)		
6	1.920	4·88 (12)	5·07 (10) 5·10 (3)	$C_{S}V_{3}O_{8}$
6	1.965	5.06(3)	5.19(3)	CoV_2O_6
6 6	1.969	5.13(1)	5.25(1)	MgV ₂ O ₆
0	1.944	4.96 (13)	5.12 (12)	CdV_2O_6

most as good a fit to the valence (5.4%) deviation as against 4.0%) as those calculated with the 54 parameters of the 27 individual curves. The universal curves thus lead to a considerable simplification in the concept of bond strength for isoelectronic series of ions. It is perhaps surprising that a single set of parameters can describe bonds ranging from almost completely ionic to largely covalent, particularly as the bond-strength model of chemical bonding has always been regarded as an ionic model. However, a prior knowledge of the covalent character of a bond is not needed in order to derive or apply the theory. On the other hand, there is a close correlation between the covalence of a bond as calculated (Pauling, 1940, p. 72) from the electronegativity difference between the terminal atoms and the mean bond strength between these atoms. Fig. 4 shows the logarithm of the covalence plotted against the logarithm of the mean cation-oxygen bond strength for 'cations' with 18, 36 and 54 electrons. The relationship between the covalence (f'_c) and the bond strength (s) for M-O bonds can be described by the empirical equation

$$f_c' = as^M \tag{11}$$

where a=0.49 v.u. and M=1.57. Similar curves can be drawn for other isoelectronic series (Fig. 5) with the values of a and M given in Table 6. It follows from equation (10) that the covalence is related also to bond length by

$$f_c' = a \left(\frac{R}{R_1}\right)^{-MN_1} \tag{12}$$

with the values of a and MN_1 given in Table 6. The constancy of the values of MN_1 for the 'ions' with 2, 10 and 18 electrons is rather striking but may be fortuitous.

Table 6. Parameters relating covalence to bond strength and bond length in the equations $f'_c = as^M = a(R/R_1)^{-MN_1}$

Number of electrons

Number of electrons			
in core	<i>a</i> (v.u.)	М	MN_1
0	0.67	1.8-2.0*	4*
2	0.60	1.73	7.04
10	0.54	1.64	7.04
18	0.49	1.57	7.04
28	0.60	1.50	9.08
36	0.49	1.57	
46	0.67	1.43	
54	0.49	1.57	

* Estimated from the trends in the value of M in subsequent series.

The relationships between the covalence and the bond strengths of M-O pairs have been calculated on the basis of mean values but it would not be unreasonable to suppose that these relationships could be used to calculate the covalence of an individual bond on the basis of its known bond strength. This would be in accord with the notion that the covalent character of a bond depends amongst other things on the oxidation state and coordination number of its terminal atoms. If these relationships can be applied to individual bonds they would allow a quantitative prediction of the difference in character between the very short and the very long V-O bonds found in V₂O₅. The short bond (R =1.586 Å) has a bond strength of 1.86 v.u. and a covalence of 1.30 v.u. giving 30% ionic character; on the other hand the long bond (R = 2.787 Å) has a bond strength of 0.103 v.u. and a covalence of 0.013 v.u. giving 87% ionic character.

Hydrogen provides another interesting example of the relationship between bond strength and covalence. Typically it forms two or more bonds between oxygen atoms in crystals. One of these is usually strong $(s \sim 0.85)$, the others are weak $(s \leq 0.15)$. The assignment of the electronegativity of hydrogen is based on a consideration of the strong hydrogen bond and leads to a covalent character of 67% for a single O-H bond. In crystals this covalence is shared very unequally between the two or more bonds. Using a value of 1.8 to 2* for the exponent *M* in equation (11) the strong bond is calculated to have about 58% and the weak bond 10% covalent character, according well with the view that the weak O-H bond is primarily electrostatic. Both in

* See Table 6. The conclusions drawn here are not very sensitive to the exact value of M chosen.

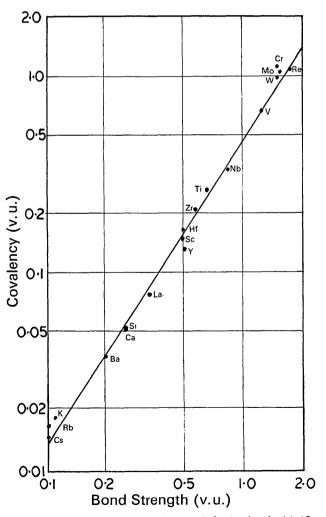


Fig. 4. M-O covalence vs. bond strength for 'cations' with 18, 36 and 54 electrons. The scales are logarithmic.

strength and character the weak O–H bond is similar to a typical Na–O bond, an observation that agrees well with the alkali metal-like behaviour of NH_4^+ .*

The fact that the relations between bond lengths, bond strengths and covalence are the same for all isoelectronic ions suggests a simple model of chemical bonding in which the electrons are divided into core electrons and valence electrons. The latter are all associated with bonds, one electron pair per unit bond strength, and are responsible for the attractive force in the bond. The core electrons give rise to the Born repulsion between atoms. Equilibrium is achieved locally when these two opposing forces are equal, the stronger the bonding force, the closer the distance of approach.

The properties of the core are described by the fitted parameters given in Table 2 and illustrated in Fig. 2. R_1 describes the sum of the radii of the cation and oxygen cores for a unit bond strength and N describes

* Khan & Baur (1972) have recently reviewed the similarities and differences between NH_4^+ and alkali metals.

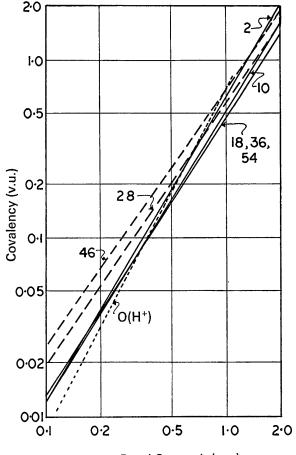




Fig. 5. M-O covalence vs. bond-strength curves for the various isoelectronic series. The numbers associated with each curve indicate the number of electrons in the 'cation' (core electrons). The scales are logarithmic.

the 'hardness' of the core, larger values representing harder cores. Both the properties of the core and the properties of the bond depend only on the numbers of electrons in each, and since these are known from classical valence theory, the prediction of bond lengths becomes remarkably simple.

This model is somewhat different from the ionic picture usually associated with a bond-strength formalism. In the ionic picture no distinction is made between core and valence electrons but the atoms are treated as negatively or positively charged spheres. In our approach all atoms are treated as positively charged spheres (nucleus plus core electrons e.g. Li⁺, S⁶⁺ and O^{2+} rather than O^{2-}) held together by the negative valence electrons. In order to apply the theory it is not necessary to know whether the electrons are placed in the bond symmetrically (covalent bond) or asymmetrically (ionic bond), but the correlation between the covalence and the bond strength suggest that the position of the electrons with respect to the oxygen atom varies in a uniform way with bond strength, residing completely on the oxygen atom at zero strength (ionic limit) and moving towards the centre of the bond for bond strengths of about two (covalent limit).

Our model differs from the Valence Shell Electron Pair Repulsion (VSEPR) theory of Gillespie & Nyholm (1957) in that we assume that the non-bonding electron pair belongs not to the valence shell but to a spherically symmetrical core. However, the VSEPR theory is most useful in discussing angles subtended by the ligands at the cation while our theory is more useful in discussing distances.

3. Distorted coordination polyhedra

The stereochemistry of many elements can be discussed quite effectively using a charged-hard-sphere model in which the properties of an atom are described solely by its ionic charge (valence) and its radius. Such atoms tend to have rather symmetric coordination with bonds of equal length. The stereochemistry of such ions is not difficult to understand in terms of most theories of chemical bonding. More recently the stereochemistry of atoms which typically have rather distorted environments (e.g. Cu²⁺, Sb³⁺ and O²⁻) have been discussed in terms of the electron configuration in such theories as the ligand field theory (Orgel, 1960) and the valenceshell electron-pair repulsion theory (Gillespie & Nyholm, 1957). In still further cases (e.g. alkali metals and V^{5+}) distorted environments occur for which no really predictive theory has yet been proposed. For atoms in such distorted environments the concept of ionic radius has proved difficult to apply and, because of the uncertainty in assigning a coordination number, it has also been difficult to use theories based on mean bond strength.

The use of R-s curves resolves this problem by avoiding the concept of coordination number altogether. In cases where the coordination is regular, the bond strength for each bond can be calculated by dividing the valency by the coordination number (Pauling, 1929; Baur, 1970), and arguments based on average bond length or bond strength lead to the same conclusions. But the method of calculating individual bond strengths can be applied just as easily in cases where the coordination is very irregular. Two examples are given in Tables 7 and 8 and are illustrated in Fig. 6 where the average bond length and the bond-strength sums are plotted as successively longer bonds are added to the coordination spheres around Na⁺ and V⁵⁺. In the bond-strength-sum graph, regular coordination appears as a straight line and irregular coordination as a curved line, but both level off at values of $\sum s = z$. The inclusion of very long bonds makes little difference to the bondstrength sums but alters the average bond length dramatically indicating that for these ions the concept of average bond length must be used with great caution.

In practice there will always be an arbitrary cut-off in the number of bonds for which strengths are calculated and this makes a small difference to the sums. At 3.0 Å the bond strength for a Na–O bond is 0.05 v.u. and inclusion of bonds in the range 3.0 to 4.0 Å corresponds to change in the value of R_0 of about 0.01

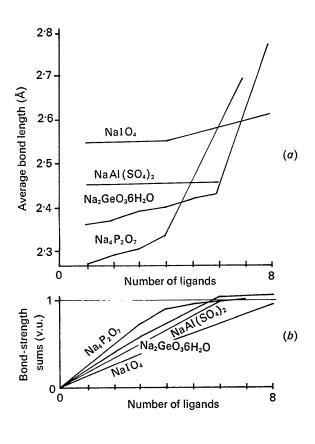


Table 7. Bond lengths in selected sodium compounds

For reference see the Appendix

Ligand number	Na ₄ P ₂ O ₇	$NaAl(SO_4)_2$	Na ₂ GeO ₃ .6H	20 NalO₄
1	2.271	2.453	2.36	2.55
2	2.314	2.453	2.37	2.55
3	2.326	2.453	2.43	2.55
4	2.441	2.453	2.45	2.55
5	2.890	2.453	2.46	2.61
6	3.156	2.453	2.52	2.61
7	3.294		3.68	2·61
8			3.85	2.61

Table 8. Bond lengths in selected vanadates

For reference see the Appendix

Ligand	Li₃VO₄	KVO ₃	V O	$Ca_{3}V_{10}O_{28}.17(H_{2}O)$
number	$L_{13}VO_4$	$\mathbf{K}\mathbf{VO}_{3}$	V 205	$(a_3 v_{10} v_{28}, 17(n_2 0))$
1	1.714	1.652	1.586	1.681
2	1.717	1.661	1.782	1.681
3	1.717	1.806	1.878	1.903
4	1.720	1.806	1.878	1.903
5			2.023	2.135
6			2.787	2.135

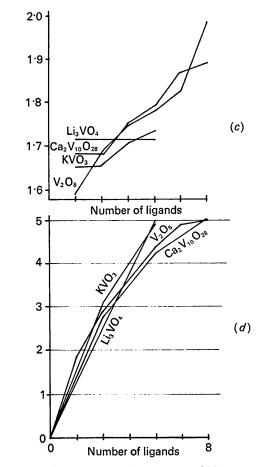


Fig. 6. (a) Average Na-O bond length as successively longer bonds are included. (b) Bond-strength sums around Na as successively longer bonds are included. (c) Average V-O bond length as successively longer bonds are included. (d) Bond-strength sums around V as successively longer bonds are included.

Å. The parameters given in Table 1 were calculated using a cut-off at around 3 Å, the point at which Nacation distances begin to appear in the coordination sphere. This cut-off point also gives the best agreement with the universal curves.

Baur (1970) is able to predict the variation of the bond lengths in distorted coordination by recognizing that where the sum of Pauling's average bond strengths around the oxygen atom is greater than the valence, the structure will compensate by weakening (lengthening) the bond to it. He obtains linear regressions between the amount of the lengthening and the excess bond strength. Although his method gives a direct prediction of the variation of individual bond lengths from the mean, it requires a knowledge of the coordination numbers in the structure and a separate curve for a given cation for each coordination number. His linear regressions are equivalent to linear R-s curves and for the examples he discusses this approximation is reasonable since the average bond length is constant from one site to another. On the other hand for strongly distorted coordination such as is found in octahedral V5+ and Cu^{2+} the average bond lengths show a significant variation. From the non-linear shape of our R-s curve we expect that such highly distorted coordination would lead to a larger average bond length. The correlation between the average bond length (\bar{R}) and the degree of distortion (Δ = mean-square relative deviation from the average) is shown in Fig. 7 for V^{5+} . Similar data for Cu²⁺, Mg²⁺ and Li⁺ are given by Shannon & Calvo (1973a, b). Linear regressions of average bond length on distortion yield the parameters given in Table 9 for V^{5+} , Cu^{2+} , Mg^{2+} , Li^+ , Zn^{2+} and Co^{2+} . For ions showing large distortion the correlation is remarkably high. For only slightly distorted octahedra other effects, such as anion coordination, become more important and this is illustrated by the better fit that is obtained in these cases when the average bond length is replaced by the effective ionic radius which takes anion coordination into account (Shannon & Prewitt, 1969).

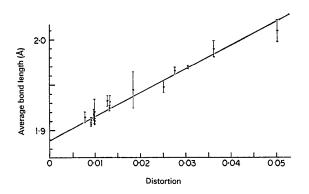


Fig. 7. Average V–O bond length as a function of distortion (mean-square relative deviation of bond length from the average) in VO₆ octahedra.

Table 9. Variation of mean M–O distance and effective ionic radius in octahedral environments as a function of distortion

The first equation given is that fitted to the observed values of \bar{R} and Δ , the second (in parentheses) is that predicted from the R-s curves using equation (13).

The effective ionic radii (\bar{r}) are calculated by subtracting the oxygen radius corrected for anion coordination from the mean bond length (Shannon & Prewitt, 1969).

$\Delta = \frac{1}{6} \sum_{i=1}^{6} [(R_i - \bar{R})/\bar{R}]^2 \text{ where } \bar{R} = \text{average bond length in Å.}$					
Atom	max ⊿	Equation	Correlation coefficient	Goodness of fit	
V ⁵⁺	0.0576	$\bar{R} = 1.887 + 2.62\Delta$	0.98	0.008	
v	0.0370	$(\bar{R} = 1.855 + 5.8\Delta)$	0 98	0 000	
Cu ²⁺	0.0316	$\bar{R} = 2.095 + 3.60\Delta$	0.89	0.016	
		$(\bar{R} = 2.084 + 6.64)$			
Mg ²⁺	0.0156	$\bar{R} = 2.094 + 8.31 \varDelta$	0.72	0.021	
		$(\bar{R}=2.098+6.3\Delta)$			
		$\bar{r} = 0.728 + 8.86\Delta$	0 ·77	0.018	
Li+	0·0 148	$\bar{R} = 2.159 + 8.42 \Delta$	0.31	0.030	
		$(\bar{R} = 2 \cdot 167 + 5 \cdot 3\varDelta)$			
		$\tilde{r} = 0.784 + 9.02\Delta$	0.79	0.035	
Zn ²⁺	0.0071	$\bar{R} = 2.099 + 7.7 \Delta$	0.64	0.021	
		$(\bar{R} = 2.111 + 6.3\Delta)$			
		$\bar{r} = 0.736 + 8.20\Delta$	0.74	0.016	
Co ²⁺	0.0046	$\vec{R} = 2.106 + 7.38\Delta$	0.42	0.019	
		$(\bar{R} = 2.118 + 6.44)$			
		$\bar{r} = 0.734 + 11.7\Delta$	0.70	0.016	

The high correlation between the average bond length and the distortion indicates that the V–O bondstrength-bond-length curve, for example, could be fitted with a quadratic equation. Such an equation can be obtained by expanding equation (5) by the binomial theorem to second order and this in turn leads to the linear relationship between average bond length and distortion:

$$\bar{R} = R_r + \frac{R_r(N+1)}{2}$$
 (13)

where R_r is the expected bond length for regular coordination. Table 9 list in parentheses the values for the coefficients of equation (13) predicted from the R-scurves. The agreement with the fitted values is satisfactory where the distortion is small but for Li, Mg and particularly Cu and V it is clear that the higher-order terms in the expansion cannot be neglected. As these depend on the nature of the distortion they will vary from one cation site to another and cannot be calculated for a general case.

4. Comparison with elastic properties

The assumption underlying the use of bond lengths to measure bond strength is that the strength of a bond can be measured by its ability to overcome the short-range repulsive force that exists between any two atoms. The parameter which measures this repulsion in an ionic theory is the Born exponent (n).*

^{*} Although the inverse power law is believed to be inferior to the exponential form (Born & Mayer, 1932; Linnett, 1940; Fumi & Tosi, 1964), Anderson (1970) has concluded that it is satisfactory for analysing the pressure derivatives of the bulk modulus of oxides.

Zachariasen (1931) has shown that in an ionic compound which exists in both the NaCl and the CsCl structure,

$$R = \sqrt[n-1]{\frac{nB}{Az_1 z_2 e^2}}$$
(14)

and

$$\frac{R_{\rm CsC1}}{R_{\rm NaC1}} = \sqrt[n-1]{\frac{nB_{\rm CsC1}}{nB_{\rm NaC1}} \cdot \frac{A_{\rm NaC1}}{A_{\rm CsC1}}}$$
(15)

where B = Born repulsion coefficient and A = Madelungconstant. Zachariasen assumed B to be proportional to the coordination number of the cation and A not to vary greatly from one structure to another. The latter assumption would appear to be fairly drastic but Templeton (1953) has shown that in a variety of structures the reduced Madelung constant ($\bar{\alpha}$ = average atomic Madelung constant) can be expressed as a function of average coordination number ($\bar{\nu}$) in the form:

$$\bar{\alpha}=1.89-\frac{1}{\bar{\nu}}.$$

If $\bar{\nu}$ varies from 4 to 6, this results in only a 6% variation in $\bar{\alpha}$. Using equation (15) Zachariasen was able to rationalize the change in interatomic distance with coordination number for a large variety of ions. Because Pauling's (1929) definition of bond strength merely involves z and v, it is apparent from equation (14) that

$$s = \frac{z}{v} = \frac{K}{R^{n-1}}$$
 (16)

The Born exponent can also be found from compressibility measurements. Anderson (1970) and Anderson & Anderson (1970) have evaluated *n* using the pressure derivative of the bulk modulus and their values (calculated using the power-law repulsion term) are compared in Table 10 with the values calculated from the R-s curves. With the exceptions of BeO, CaO, SiO₂ and TiO₂ there is satisfactory agreement. Anderson & Anderson have found the elastic behaviour of

 SiO_2 and TiO_2 anomalous and attribute the effect to bond bending and directional forces. The values of ndetermined from the elastic constant for BeO and for CaO (Son & Bartels, 1972) differ by a factor of two from the values we have obtained. Fisher (1973) has suggested that the origin of this discrepancy may be in the derivation of *n* from the pressure derivative of the bulk moduli. Using the relationships derived by Anderson & Liebermann (1970), Fisher obtains values of nfor single-crystal MgO and CaO from the elastic modulus, C_{44} , and the bulk modulus, K. These values agree well with the values of n determined from crystallographic data in Table 10. It would be interesting to see if the discrepancy between the two values of *n* for BeO could be removed in a similar manner. In the light of the rather drastic assumptions implicit in equation (16), the agreement between the Born exponent determined from the elasticities and that determined from the crystal structures is encouraging.

Uses of bond-strength-bond-length curves

Previous workers have used bond-strength-bondlength curves for a number of practical purposes related to the evaluation of crystal-structure determinations (see Donnay, 1972). Our curves can achieve the same results with considerably more generality and flexibility as is illustrated in the examples given below.

1. Testing for the correctness of a structure

Bond-strength curves provide a convenient method of testing the plausibility of a proposed crystal structure, since the bond-strength sums around any atom should normally lie within 5% of the valence. Larger discrepancies in well refined structures might indicate an incorrect structure. The effect is well illustrated by the structure of $Zn_3(BO_3)_2$ which was originally refined by Garcia-Blanco & Fayos (1968) in the space group *Ic*. Baur & Tillmans (1970) subsequently refined it in the space group I2/c and although they obtained the same

Table 10. The Born exponent n evaluated from elastic constants and from bond-strength parameters

Compound	From elastic constants	Born exponent From individual bond-strength curves§	From universal bond-strength curves§
BeO	9.5*	5.3	5·0 [°]
MgO	6.5,* 4.7†	6.0)
Al_2O_3	5.0*	6.0	
α-SiO2 SiO2(stishovite)	12·2* 14·0*	5.5	5.3
MgAl ₂ O ₄	5.6*	6.0	
Mg2SiO4	6.5*	5.7	j
CaO TiO₂(rutile)	8·7,* 11·0,† 6·11 13·3*	‡ 7·0 5·0	} 5.5
Fe ₂ O ₃	6.6*	6.3	6.1
ZnO	7.4*	6.0	7.0
† So ‡ Fi	nderson (1970) and Anderson n & Bartels (1972) sher (1973) : N + 1	& Anderson (1970))

agreement index (R=0.12) as Garcia-Blanco & Fayos they preferred their structure on chemical as well as crystallographic grounds (Baur, 1970). The bondstrength calculations confirm Baur & Tillmans's preference, giving a mean deviation of the bond-strength sums from the valence of 13% for the refinement in *Ic* and only 3.4% for that in I2/c (see Table 11).

Table 11.	Bond-strength	sums in	two	structures
	proposed for	$Zn_3(BC)$	$)_{3})_{2}$	

Space group	z	<i>Ic</i> Garcia-Blanco & Fayos (1968)		<i>12/c</i> Baur & Tillmans (1970)
	-		• •	. ,
Zn(1)	2	1.99	2.07	1.95
Zn(2)	2	2.24	1.69	1.96
Zn(3)	2	2.08	1.80	1.99
B (1)	3	3.23	3.44	3.12
B(2)	3	2 ·76	3.67	3.01
O(1)	2	2.39	2.11	2.05
O(2)	2	2.08	1.66	1.95
O(3)	2	1.97	2.55	2.12
O(4)	2	2.31	1.96	2.08
O(5)	2	1.82	1.95	1.91
O(6)	2	1.77	2.39	1.92
Mean deviation from z		13	0/0	3.4 %
		15	/0	5470

2. Determining site occupancy in solid solutions

When two atoms of different valence but belonging to the same or neighbouring elements (e.g. Al^{3+} and Si^{4+} ; Fe^{2+} and Fe^{3+}) are present in a crystal, it is usually not possible to determine directly from X-ray diffraction studies the ratio of each on any particular site. The occupancy can be inferred from a study of the bond lengths between the site and its first neighbours since the lower valence ion gives rise to longer bonds (see for example Evans, 1960; Smith & Bailey, 1963; Ribbe & Gibbs, 1969). Such methods depend on the construction of bond-length-site-occupancy curves for each pair of ions. The use of bond strength provides a more general way of assigning site occupancy since such disordered atoms will normally have the same or very similar R-s curves, and the calculated bond-strength sums will be equal to the average valence of the site. Two examples are given. In Table 12 the distribution of Al³⁺ and Si⁴⁺ in sanidine, (Na, K)AlSi₃O₈, is predicted using the universal bond-strength curve, the bond-strength curve corrected for oxygen coordination, and ionic radii (Shannon & Prewitt, 1969). These are compared with the site occupancy determined from neutron diffraction and the agreement is quite satisfactory. The method can be used with less precision to calculate the ratio of K to Na on the alkali metal site.

Table 13 compares the bond strengths in three ferrospinels. The bond-strength sums at each site have been calculated with the R-s curves appropriate to each species present and the weighted average bond strengths are shown for three cases–(a) normal spinel, (b) inverse spinel and (c) random spinel. In each case these are compared with the expected average valence. Table 13 clearly indicates that Fe₂TiO₄ is an inverse spinel and Fe₂ZnO₄ is normal in agreement with assignments given by other workers. The result for Fe₃O₄ is not quite as unambiguous but suggests that the Fe²⁺ ions are disordered over all sites, a result at variance with accepted view that Fe₃O₄ is an inverse spinel.

3. Prediction of positions of hydrogen atoms

The R-s curve for hydrogen has been calculated using hydrogen positions determined by neutron dif-

Table 12. Cation distribution in sanidine

Structural information from Prewitt (1973) Estimated errors given in parentheses.

Site	Bond strength	Universal	Bond strength	Ionic radii	Neutron
	sums	bond-strength	corrected for	(Shannon &	diffraction
	(universal curves)	sums	O coordination	Prewitt, 1969)	(Prewitt, 1973)
T1 T2 K, Na	3·72 3·87 1·13 (as K) 0·79 (as Na)	72 (10)% Si 87 (10)% Si 62 (30)% K 38 (30)% Na	73 % Si 82 % Si	71 %Si 92 % Si	68 % Si 82 % Si 88 % K 12 % Na

Table 13. Bond strengths in Fe_2XO_4 spinels

			$P(Fe^{3+})$	$P(\mathrm{Fe}^{2+})$	$P(\mathbf{X})$	P(nor	rmal)	P(inv	erse)	P(ran	dom)
Х	Site	$R(\text{\AA})$	Table 4-2	Table 4-1	Table 4-2	Theor.	Obs.	Theor.	Obs.	Theor.	Obs.
Ti	Α	2.00		2.01	2.49	4.00	2.49	2.00	2.01	2.67	2 ·17
	В	2.03		2.78	3.49	2.00	2 ·78	3.00	3.16	2.67	3.02
Zn	A	1.97	2.25		1.93	2.00	1.93	3.00	2.25	2.67	2.14
	В	2.02	2.96		2.48	3.00	2.96	2.50	2.72	2.67	2.80
Fe	A	1.887	2.80	2.59		2.00	2.59	3.00	2.80	2.67	2.73
	В	2·0 56	2.69	2 ·76		3.00	2.69	2.50	2.72	2.67	2 ·71

References

Fe2TiO4 Ishikawa, Sato & Syono (1971).

 Fe_2ZnO_4 Chalyi & Rozhenko (1956); Bałanda, Szyłula, Dimitrijević & Todovović (1969). Fe_3O_4 Hamilton (1958) fraction. In many cases it is not possible to detect hydrogen atoms using X-ray diffraction and where it is possible their positions are not only less accurately known but differ systematically from the positions found by neutron diffraction. Consequently the bond strength curves for H–O bonds quoted in Table 1 are only appropriate for bonds determined by neutron diffraction.

Often it is of interest to know where the hydrogen atoms are to be found in the crystal in those cases where their presence is not detectable by X-ray diffraction. Frequently the presence of short O-O distances is indicative of a hydrogen bond and the hydrogen atoms can be assigned qualitatively without much ambiguity. Donnay & Allman (1970) have shown that the calculation of bond-strength sums provides additional information about the location of hydrogen atoms. Our bond strengths can be used in a similar way but have the additional advantage of ease in calculation. Table 14 shows the calculations of bond strengths for

MgSO₄.4H₂O both with and without the inclusion of the hydrogen atoms. Table 15 shows similar calculations for Na₂CO₃.H₂O where the hydrogen positions are not known. Estimates of O-H distances based on predicted bond strengths give values for the O-H···O distances that are in striking agreement with those observed. A somewhat different prediction for the positions of the hydrogen atoms based on different considerations has recently been made by Baur (1972).

4. Calculation of effective ionic radii

The R-s functions given in Table 3 allow one to obtain more statistically valid effective ionic radii than were previously possible and in addition allow extrapolation to obtain radii for which no examples have yet been found. Some examples of previously undetermined radii and of radii differing significantly from values in Shannon & Prewitt (1969) are listed in Table 16. Table 15. Prediction of hydrogen atom positions in $Na_2CO_3 \cdot H_2O$ using universal bond-strength curves

Structural information from Dickens, Mauer & Brown (1970).

	Bond streng sums witho H	gth stre	ated bond ngths for drogen H(2)	Estimated bond strength sums with H
Na(1) Na(2) C	1·02 1·11 3·98			1·02 1·11 3·98
O(1) O(2) O(3)	1·88 1·98 1·80	0.16	0.24	2·04 1·98 2·04
O(4) Hydrogen	0.45 bond	0·84 1·00	0·76 1·00	2.05
	Bond strength	Bond length (Å)	Predicted O-H····	
O(4)-H(1) H(1)-O(1)	0.16	0·93 2·1	} 3.0	2.907 (2)
O(4)-H(2) H(2)-O(3)		0∙96 1∙7	} 2.7	2.684 (2)

Table 16.	Effective	ionic	radii	of	ions	in	unusual
	co	ordin	ation	5			

Ion	CN	Radius
Li+	v	0.69
B ^{3 +}	VI	0.22
P ⁵⁺	VI	0.38
S ^{6 +}	VI	0.29
V ⁵⁺	V	0.43
V ⁵⁺	VI	0.20
Mn ² +	IV	0.65
Mn ²⁺	v	0.74
Fe ² +	v	0.70
Fe ² +	VIII	0.91
Fe ³⁺	V	0.56
Ge ⁴⁺	v	0.46
As ^{5 +} As ^{5 +}	V	0.40
As ⁵ +	VI	0.46

Table 14	Rond strengths in	n MoSO.4H.O	calculated using	universal curves
1 4010 1 4. 3	Dona bacing and a	111500411120	curculated ability	Millerou. chieco

			Structura	l informa	tion from	n Baur (19	964)			
	Bond strength sums without H	H(1 <i>A</i>)	H(1 <i>B</i>)	Bond str $H(2A)$	engths fo H(2B)	r hydroge $H(3A)$	en atoms H(3B)	H(4 <i>A</i>)	H(4 <i>B</i>)	Bond strength sums with H
Mg S	2·14 6·08		. ,							2·14 6·08
O(1)	1.84						0.16			2.00
O(2)	1.89				0·07 0·09					2.05
O(3)	1.50	0.15	0.17	0.14		0.10		0.15	0.19	2.00
O(4)	1.55	0.00	0.96	0.14		0.12		0.12		1·97 2·05
O(W1) O(W2)	0·37 0·35	0.82	0.86	0.82	0.90	0.86	0.78			2·03 2·07 1·99
O(W3) O(W4)	0·36 0·36					0.90	0.78	0 ·84	0.79	1.99
Bond stre sums for O-H····O	н	0.97	1.03	0.96	1.06	0.98	0.94	0.99	0.98	
distance		2· 884	2.754	2.835	3·282 3·042	2.860	2.833	2.831	2.734	

5. Testing for positional disorder or anomalous thermal motion

As mentioned above, thermal motion tends to decrease the apparent bond length as determined by diffraction methods. In the present discussion no allowance has been made for this effect on the assumption that the extent of the shortening will be approximately the same for all bonds. If, however, the thermal motion is anomalous large (or small) the apparent bond lengths will be shortened (or lengthened) from the values used in this study, giving rise to anomalously high (or low) bond strengths. The effect is well illustrated by the bond-strength sums calculated around the bridging oxygen atoms in pyrophosphates (Table 17). All the sums are somewhat high but they tend to increase as the P-O-P angle increases reaching values between 2.3 and 2.4 in those structures where the P-O-P angle is constrained by symmetry to be 180°. In these latter cases the thermal parameters of the bridging oxygen atom are large, indicating that the atom is normally disordered away from the centrosymmetric position. The instantaneous P-O-P angle is thus less than 180° and the corresponding P-O distances longer than the values of 1.56 Å normally found in these crystals. The large calculated bond-strength sums in this case are thus indicative of either a positional disorder or an anomalous thermal effect.

Table 17. Bond-strength sums around the bridging oxygen atoms in pyrophosphate groups

	Bond-strength sum	P-O-P	
	(Table 4-2)	(°)	Reference [†]
$Cd_2P_2O_7$	1.99 (5)	132	69 CJCHA 47 3409
$Na_4P_2O_7$	2.10(1)	128	72 CJCHA 50 2519
α -Zn ₂ P ₂ O ₇	2.10(2)	139	70 JSSCB 1 120
α -Sr ₂ P ₂ O ₇	2.12(5)	131	68 ACSAA 22 1419
β -Ca ₂ P ₂ O ₇	2.12(4)	131	66 ACCRA 21 942
α -Ca ₂ P ₂ O ₇	2.14(3)	130	68 INOCA 7 1345
α -Mg ₂ P ₂ O ₇	2.18(4)	144	67 ACCRA 23 289
β -Ca ₂ P ₂ O ₇	2.22(4)	138	66 ACCRA 21 942
α -Zn ₂ P ₂ O ₇	2.23(3)	148	70 JSSCB 1 120
α -Cu ₂ P ₂ O ₇	2·26 (3)	157	67 ACCRA 22 665
SiP ₂ O ₇ -III	2.27(4)	139	70 ACBCA 26 233
β -Zn ₂ P ₂ O ₇	2.31(2)	180	65 CJCHA 43 1147
β -Mg ₂ P ₂ O ₇	2.38 (2)	180	65 CJCHA 43 1139

 \dagger The references are given in the same form as in the Appendix.

Table 18. Structures used in the evaluation of bond-strength-bond-lengths parameters

H11+) 69 ACBCA 25 311 BE (H2,0) 4 5.04 69 ACBCA 25 1604 BC (H2,0) 4 5.04 66 ACCRA 25 1604 BC (H0,3) 2 (D2.0) 4 66 ACCRA 22 162 NA AL (S.04) 2 (H2.0) 12 66 ACCRA 22 162 NA AL (S.04) 2 (H2.0) 12 66 ACCRA 21 383 CA AL (S.04) 2 (H2.0) 9 66 ACCRA 21 383 CA AL (S.04) 2 (H2.0) 9 66 ACCRA 21 365 NF (BS.03) (H2.0) 9 66 ACCRA 20 965 H 1 103 66 ACCRA 21 980 CA HA (I AL (S.04) 2 (H2.0) 9 66 ACCRA 21 980 CA HA (I AL (S.04) 2 (H2.0) 9 66 ACCRA 21 980 CA HA (I AL (S.04) 2 (H2.0) 9 66 ACCRA 21 980 CA HA (I AL (S.04) 2 (H2.0) 9 67 ACCRA 19 396 CA HA (I AL (S.04) 2 (H2.0) 9 66 ACCRA 19 396 CA HA (I AL (S.04) 2 (H2.0) 9 66 ACCRA 19 396 CA HA (I AL (S.04) 2 (H2.0) 9 66 ACCRA 19 396 CA HA (I AL (S.04) 2 (H2.0) 9 66 ACCRA 19 396 CA HA (I AL (S.04) 2 (H2.0) 9 66 ACCRA 19 396 CA HA (I AL (S.04) 2 (H2.0) 9 66 ACCRA 19 396 CA HA (I AL (S.04) 2 (H2.0) 9 66 ACCRA 19 196 CA HA (I AL (S.04) 2 (H2.0) 9 66 ACCRA 19 196 CA HA (I AL (S.04) 2 (H2.0) 9 66 ACCRA 19 196 CA HA (I AL (S.04) 1 (H2.0) 9 66 ACCRA 19 196 CA HA (I AL (S.04) 1 (H2.0) 9 66 ACCRA 19 197 CA (S.04) 1 (H2.0)	67 ACCRA 22 37 NA2 SI 03 SI2 05 68 ACBCA 24 117 BETA-MA2 SI2 05 71 ACBCA 27 2124 NA3 P.3 (H2.0)12 67 ACCRA 12 72124 NA3 P.3 (H2.0)11 67 ACCRA 12 72124 NA3 P.3 (H2.0)11 67 ACCRA 12 72124 NA3 P.3 (H2.0)11 70 ACBCA 25 107 NA5 P.3 08 (H2.0)14 70 ACBCA 25 107 NA5 P.3 08 (H2.0)14 70 ACBCA 25 107 NA5 P.2 C5 0.02 71 ACBCA 27 126 NA3 P.3 08 (H2.0)14 70 ACBCA 27 126 NA3 P.3 08 (H2.0)14 70 ACBCA 27 126 NA3 P.3 08 (H2.0)12 70 ACSAA 27 126 NA3 P.3 08 (H2.0)12 70 ACSAA 27 126 NA3 P.3 08 (H2.0)12 70 ACSAA 22 1357 NA4 CR3 08 07H 71 ACBCA 27 2126 NA3 P.3 08 (H2.0)12 70 ACSAA 21 1261 NA4 CR3 08 07H 71 ACBCA 26 159 NA2 CR3 08 07H 70 ACSAA 21 1261 NA4 CR3 08 07H 71 ACBCA 26 1597 NA4 CR3 08 07H 70 ACSAA 21 1261 NA4 CR3 08 07H 70 ACSAA 21 1261 NA4 CR3 06 (H2.0)12 70 ACSAA 21 1261 NA4 SK4 0210 035 (0.H14 71 ACBCA 26 1597 NA4 SK3 05 (H2.0)12 70 ACSAA 21 1261 NA4 SK4 0210 035 (0.H14 71 ACBCA 26 1596 NA2 H 450 04 (H2.0)7 70 ACSAA 21 1261 NA4 SK4 0210 035 (0.H14 71 ACBCA 26 1596 NA2 H 450 04 (H2.0)7 70 ACSAA 21 1261 NA4 SK4 02 J0 03 (H2.0)12 70 ACSAA 21 1261 NA4 SK4 02 J0 (H2.0)7 70 ACSAA 21 1267 NA4 SK4 02 J0 (H2.0)7 70 ACSAA 21 1267 NA4 SK4 02 J0 (H2.0)7 70 ACSAA 21 1267 NA4 SK4 02 J0 (H2.0)7 70 ACSAA 21 1268 NA2 H 42 J0 (H2.0)7 70 ACSAA 21 1268 HC AL 20 (H2.0)7 70 ACSAA 21 1268 HC AL 20 (H2.0)7 70 ACSAA 23 1468 HC AL 20 (H2.0)7 60 AMHIA 50 00 HC AL 20 SI 00 A 60 AMHIA 50 00 HC AL 20 SI 00 A 60 AMHIA 50 00 HC AL 20 SI 00 A 60 AMHIA 50 00 HC AL 20 SI 00 A 60 AMHIA 50 00 HC AL 20 SI 00 A 60 AMHIA 50 00 HC AL 20 SI 00 A 60 AMHIA 50 00 HC AL 20 SI 00 A 60 AMHIA 50 00 HC AL 20 SI 00 A 60 AMHIA 50 00 HC AL 20 SI 00 A 60 AMHIA 50 00 HC AL 20 SI 00 A 60 AMHIA 50 00 HC AL 20 SI 00 A 60 AMH
07 ACCERA 22 354 42 BE-54 B(3+) 70 ACCERA 26 906 B2 03 I 63 ACCEA 24 6059 BE F10+ 0102 54 ACCERA 16 330 GAMMA-4B 022 54 ACCERA 16 330 GAMMA-4B 022 54 ACCERA 16 330 GAMMA-4B 022 54 ACCERA 16 1233 NA 8 (3,4) 64 JCFSA 4 3304 GAMAL-11 8.02 71 ACCEA 27 904 L13 B0 03 65 ACCERA 16 1233 NA 8 (3,4) 65 ACCERA 16 1233 NA 8 (3,4) 65 ACCERA 16 1233 NA 8 (3,4) 65 ACCERA 16 372 K2 B0 03 70 ACCERA 22 130 K4 GAMAL 65 MIXAS 1965 196 MC AL 504 65 ACCERA 16 372 K2 B0 03 65 ACCERA 16 370 K2 B0 03 65 ACCERA 16 350 CA 82 34 (171) 66 ACCERA 25 955 CA 82 34 (171) 66 ACCERA 25 955 CA 82 34 (171) 66 ACCERA 25 955 CA 82 34 (171) 67 ACCERA 27 677 CU 82 34 63 ACCERA 16 1237 NA 8 SI3 04 64 ACCERA 16 125 186 AC 84 05 (0,4) 65 ACCERA 16 350 CA 82 35 (0,4) 66 ACCERA 16 350 CA 82 35 (0,4) 67 ACCERA 27 905 NA 88 SI3 04 63 ACCERA 16 1237 NA 88 SI3 04 64 ACCERA 16 1237 NA 88 SI3 04 65 ACCERA 22 905 NA 4L 15 04 65 ACCERA 22 905 NA 4A L 50 04 65 ACCERA 22 905 NA 4A L 50 04 65 ACCERA 22 120 NA 4A L 50 04 65	65 ACCRA 19 396 GAMHA'LI AL O2 68 ZEKCA 126 46 LI AL 312 36 63 AHMIA 48 804 BE AL2 34 65 MIASA 1995 196 AL MG 3.04 68 ACCRA 24 158 MG AL3 312 06 66 ACCRA 24 158 MG AL3 312 06 67 ACCRA 25 182 NA AL 351 06 67 ACCRA 25 182 NA AL 351 06 67 ACCRA 151 314 NA CA AL 512 07 65 AHMIA 50 2023 MG3 AL2 513 012 61 ZEKCA 115 259 *ANOALJSITE* AL2 SI 05 63 ZEKCA 118 337 *KIANITE* AL2 SI 05 63 ZEKCA 123 81 CA3 AL2 SI 05 64 ZEKCA 123 81 CA3 AL2 SI 05 65 ZEKCA 123 151 AL P634 OS 66 ZEKCA 123 151 AL P634 OS 66 AL2 51 05 67 ZEKCA 123 151 AL P634 OS 68 AMMIA 53 1056 AL2 53 012 66 ZEKCA 123 151 AL P634 OS 66 AMMIA 53 1056 AL2 53 027

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APPENDIX

Structures and their references used in the evaluation of the bond-strength-bond-length parameters given in Tables 1, 2 and 3 are given in Table 18. The reference is given in the following form: year, journal (a list of codes used is given at the end), volume and page number. For papers marked with an asterisk, see Table 4.

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Table 18 (cont.)

TI(4+) 71 JCPSA 77 4708 71 JCPSA 55 3206 69 ZKKGA 129 222 69 ZKKGA 129 438 64 ACCRA 127 240 68 ACGRA 17 240 68 ACGRA 131 276 66 CHOAA 262 962 65 CJCHA 47 971 69 CJCHA 47 971 69 CJCHA 47 971	*ANASTASE * 11 02 *RUTIL* 11 02 BA TI 313 09 BA TI 0 SI2 07 CO 11 32 Y2 12 07 KA TI 05 CO 15 05 Y2 12 07 KA TI 05 CO 15 05 CO	62 JINCA 69 ZEKGA 69 ZEACA 70 PEP11 68 Annia 70 Annia 62 ZEKG 65 JCPSA	24 237 129 427 130 1399 358 1261 533 2733 557 20333 43 2533	MN CR2 04 MN 5E 2 GE3 012 MN 5E 23 ETA MV2 GE404 HN7 (0, H) 8 AS2 08 MN9 (0, H) 8 AS2 08 MN M0 34 MN M0 34
60 JCDSCA 32 1515 69 ACGCA 32 1514 72 CCRSC1 1645 70 ACGBCA 21 1645 66 ACCCRA 21 974 55 BCSJA 22 955 44 1FSSOA 37 393	TI LR4	66 ANMIA 69 MSAPA 67 ACCRA 69 MSAPA 69 ACBCA 71 AMMIA 68 Zekga 70 Algca	51 123 22 761 22 15247 550 19297 550 19199 557 1265 1265 1465	$ \begin{array}{llllllllllllllllllllllllllllllllllll$
61 ZEKGA 115 110 UNPU1 REF1 54 UNPU1 REF1 66 UNPA 5 2137 66 UNPA 5 2137 71 ACBCA 114 2237 71 ACBCA 114 2237 71 ACBCA 114 237 71 UNPIA 27 REF2 71 UNPIA 27 REF2	$\begin{array}{c} 1 & 0 \\ 1 & 1 & 1 \\ 1 & 1 & 1 \\ 1 & 1 & 1 \\ 1 & 1 &$	44 A 4 A 94 A 44 A 44 A 44 100 JUGANACULUU 107 RUNGULUUUUU 107 RUNGULUUUUU 107 RUNGULUUUUU 107 7 67 6 67 7 7 7 7 7	2 2 25 51 27 17 4 900000000 2 2 5 51 19 21 5 17 4 900000000 1 2 5 55 55 5 5 5 5 5 5 5 5 5 5 5 5 5 5	GAS CLEFPLESOLD HANS FEE GES 012 FE V 00 FE V 00 FE 35 FE 25 FE 25
71 CJCULA 49 (%62) JJSKCA 114 827 607 ZKKCA 114 827 607 ZKKCA 114 727 607 JNOCA 5 967 70 CHOCA 270 252 60 LANCCA 270 252 60 LANCCA 21 837 72 JSSCB 4 23 72 JSSCB 4 2	K V.03 K V.03 K V.03 K 2 N2 V.0 CA V.34 CL CA V.34 CL CA V.34 CL CA V.35 CA V.	70040000000000000000000000000000000000	2666220088882266622008882266622008888226666220098822666622009711199711997119971	LB FE 33 PV FE 03 PV FE 03 PV FE 03 PB FE 03 PB FE 03 PS FE 03 PS FE 03 PS FE 03 PS FE 03 PS FE 03 PS 012 G03 FE 012 G03 FE 012 G03 FE 012 G03 FE 012
UNPU1 REF3 UNPU1 REF3 UNPU1 REF3 71 CJCHA 49 3156 54 ACCRA 49 2156 54 BCCFA 3 736 69 BCCFA 45 2297 69 ACBC44 24 62297 69 ACBC44 24 62297	CÕ V2 06 CO2 V2 07 CO3 V2 08 ZN2 V2 07 ALPNA-N3 V2 08 CO V2 06 CO V2 06 CO V3 06 CO V3 06 CO V3 06 CO V4 06 V V0 V V0 V V0 V V0 V V0 V V0 V V0 V	54 ACCXA 70 SSC0A 65 ACCRA 65 ACCRA 66 NJHMA 67 CP SA 66 NJHMA 69 NJHMA 67 CIWYA 67 CIWYA	$\begin{array}{c} 19 & 9745 \\ 19 & 1971 \\ 19 & 9759 \\ 1966 & 430 \\ 1969 & 430 \\ 1969 & 430 \\ 1969 & 430 \\ 15 & 2885 \\ 67 & 2885 \\ 18 & 807 \\ 156 & 1399 \\ 661 & 290 \end{array}$	YE3 FES 015 0/3 FES 012 H03 FES 012 H03 FES 012 LU3 FES 012 LU3 FES 012 FE 0 H4 FE2 AL SI4 024 FE 0 H4 FE2 AL SI4 024 FE 0 FE 0 FE 0 FE 0 FE 0 FE 0 FE 0 FE 0 FE 0 H4 FE2 AL SI4 024 FE 0 FE
66 INOCA 5 1808 CR(6) * 70 ACSAC 26 222 70 ACSAC 26 222 70 ACSAC 26 100 * 68 CJCHA 55 1103 * 68 CJCHA 55 1103 * 68 CJCHA 92 264 69 BUFCA 92 264 69 BUFCA 92 1697 * 71 ACSAC 25 * 71 ACSAC 25	CS V3 38 (PM3 SI)2 CR.04 IN PBID F2 (CR.04) K2 (R2.07 CG (CR.04) N (CR.04) N (CR.04) R32 (CR.07) R32 (CR.07) R33 (CR.07) CR.07) CR.07	65 ACCRA 68 AHMIA 69 SCIEA 69 SCIEA 69 CIWYA 62 ACCRA 67 ACCRA 71 JUPSA 67 ZEKGA	166 1399 68 290 93 190 15 815 22 775 31 452 369 300 124 192	FË \$.04 ⁻ (H2:0)4 FE (N,H4)2 (\$.04)2 (H2:0)6 FE2TI 04 FE V2 04 FE H 04
★ 68 GUCHA 46 333 ★ 68 BUFCA 92 7687 ★ 69 BUFCA 92 7687 ★ 71 ACSAA 25 1969 ★ 71 ACSAA 25 35 ★ 71 ACCAA 13 325 ★ 70 NJMKA 133 1 ★ 10 NC2+) ★ 10 NC2+ ★ 10 NC2	RB2 CR2:07 VIII RB2 CR2:07 X AC2 CR2:04 AC2 CR2:04 HT N0:04 HT N0:04 HT N0:04 HT N0:04 CA HN SI2 C6 CA HN SI2 C6 HA C0 HA C0	50 NATUA 68 CJCHA 71 ACBCA 70 PEPI1 62 ACCRA 64 ACCRA 64 ACCRA 64 ACCRA	54 1621 22 775 17 240	
70 726464 1132 67 726464 152 717 67 726464 152 717 67 726664 154 145 67 726664 154 145 67 726664 117 132 62 72664 117 131 62 72664 117 131 62 72664 117 131 63 72664 117 131 65 72664 117 131 75 72664 117 117 117 117 117 117 117 117 117 11	2N2 MN (0.4)2 S1.04 CA2 MN (0.4)2 S1.04 * 952CLAUELTER FR F0.04)2 (P.04)2 (H2. * 952CLAUELTER FR F0.04)2 (0.4)2 (H2. MN S.04 MN S10 06 MN 10 30 MN V2 07	UNPU1 UNPU1 UNPU1 H2.0 68 ZEKGA 70 CJCHA 70 CJCHA 68 ZAACA 65 ACCRA	REF3	CO2 V2 07 CO V2 05 CO V2 05 CO3 C2 03 CO3 C2 03 CO3 C3 C3 CO3 C3 CO3 C3 CO3 C3 CO3 C3 C0 C3 C0 C3 C0 C3 C0 C3 C0 C3 C0 C3 C0 C3 C3 C3 C3 C3 C3 C3 C3 C3 C3 C3 C3 C3 C

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Table 18 (cont.)

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